BACTERIAL DIVISION

Mechanical crack propagation drives millisecond daughter cell separation in *Staphylococcus aureus*

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When *Staphylococcus aureus* undergoes cytokinesis, it builds a septum, generating two hemispherical daughters whose cell walls are only connected via a narrow peripheral ring. We found that resolution of this ring occurred within milliseconds ("popping"), without detectable changes in cell volume. The likelihood of popping depended on cell-wall stress, and the separating cells split open asymmetrically, leaving the daughters connected by a hinge. An elastostatic model of the wall indicated high circumferential stress in the peripheral ring before popping. Last, we observed small perforations in the peripheral ring that are likely initial points of mechanical failure. Thus, the ultrafast daughter cell separation in *S. aureus* appears to be driven by accumulation of stress in the peripheral ring and exhibits hallmarks of mechanical crack propagation.

ost bacteria propagate through binary fission, a process that is highly coordinated and tightly controlled to pass on genetic material equally to the two daughter cells and to regulate cell size and shape. Much of our knowledge of bacterial cell division comes from rod-shaped bacteria, which double their cell length before cytokinesis (1, 2): relatively less is known about cell division in bacteria with other shapes. Staphylococcus aureus, a model system for round bacteria, is a Gram-positive pathogen well recognized for its virulence and antibiotic resistance (3, 4). To divide, S. aureus builds a septum, generating two hemispherical daughter cells (5, 6). After construction, the septal wall exists as two flat, parallel plates, and the walls of the two daughter cells are connected only through a narrow peripheral ring (Fig. 1A) (7). Presumably, resolution of this peripheral wall ring leads to daughter cell separation, which is accompanied by a shape conversion of the daughter cells from hemispheres to spheres. This shape change has previously been assumed to occur through expansion of the septum to twice its original surface area, which would double the cell volume (5, 8). It remains unclear how exactly the peripheral ring is resolved to allow the daughter cells to separate, particularly given that the S. aureus cell wall is quite thick (20 to 30 nm) (9).

Previous video microscopy-based observations of *S. aureus* cell division have described daughter

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cell separation as a dramatic "popping" event with no detectable intermediate stages (10, 11). To address the time scale and mechanism of the popping, we used phase contrast microscopy with a temporal resolution of 1 ms. At this frame rate, we occasionally observed intermediate stages of popping, whereas most separations occurred within one or two frames (<2 ms, or 1/10⁶ of the cell cycle) (Fig. 1B and movie S1). This rapid separation of the daughter cells with drastic shape change contrasts sharply with the gradual morphological changes commonly associated with cell division in other bacteria (12), suggesting that daughter cell separation in S. aureus must not be solely dependent on enzyme-mediated cell-wall remodeling. Rather, the millisecond daughtercell separation suggests the involvement of mechanical forces (13). One possibility is that the peripheral ring connecting the two daughter cells is under substantial mechanical stress before separation, so that if a random point on this ring were to fail, a crack would propagate around the periphery, separating the daughter cells but leaving them connected by a hinge point roughly opposite the position of the initial point of material failure. The observed rate of daughter-cell separation (~1 µm/ms) is well within the range of crack propagation speeds for soft biological materials (14. 15).

An essential feature of mechanical crack propagation is its dependence on stress in the material (16, 17). The primary source of stress in the bacterial cell wall is turgor pressure (18). Indeed, removing turgor pressure in *S. aureus* by disrupting the cytoplasmic membrane resulted in an average decrease in cell volume of >20% (fig. S1 and movie S2), indicating that the cell wall is normally under substantial mechanical stress. If cell-wall stress and consequent mechanical failure are contributing factors to the ultrafast cell separation, then altering turgor pressure should influence the likelihood of separation. To test this hypothesis, we exposed an unsynchronized, growing population of cells to oscillatory changes in medium osmolarity over a range of 100 to 500 mM in order to modulate turgor pressure and cell-wall stress and recorded the time of popping with respect to the phase of the oscillatory cycle for hundreds of individual popping events. We observed a large dose-dependent enrichment of popping events during the intervals when medium osmolarity was being lowered (downshift), which corresponds to an increase in turgor pressure and cell-wall stress, and a depletion of popping events during the intervals when medium osmolarity was being raised (upshift) (Fig. 1C and fig. S2). Thus, an externally induced increase in cell-wall stress promotes popping, whereas a decrease in wall stress delays popping, confirming the involvement of cell-wall stress in determining the likelihood of popping.

A further prediction of the stress-driven crack propagation model in which failure is initiated at one random point along the periphery is that after splitting, when stress has been released, the two daughter cells will remain connected at a hinge point opposite the initial site of failure. To probe the relative orientation of the two daughter cells after popping, we tracked the fate of the outer wall (Fig. 1A) relative to the septal wall after cell separation using fluorescent wheat germ agglutinin (WGA) and three-dimensional (3D) structured illumination microscopy (3D SIM). WGA binds to N-acetylglucosamine residues in the cell wall (19) and does not penetrate into the septum because of its size and can therefore be used to selectively label the S. aureus outer wall (20). In nearly all of the daughter cell pairs observed (39 of 40), the two sections of the previous outer wall were still partially connected after cell separation, and in all cases (40 of 40) they appeared to have rotated around a hinge (Fig. 1D, 10 min, and movie S3). In addition, we followed WGA-labeled live cells with epifluorescence microscopy and observed two WGA labeling patterns after popping: the hinged pattern as observed with 3D SIM (Fig. 1E, left) and a nonhinged pattern (Fig. 1E, right) that resembles the labeling pattern reported previously (20). By correlating epifluorescence microscopy to scanning electron microscopy (SEM), we realized that the nonhinged pattern corresponds to cells with their hinge points oriented at the top or bottom surface of the cells relative to the coverslip (Fig. 1F). Thus, daughter-cell separation in S. aureus is achieved through mechanical crack propagation that initiates at some point around the peripheral ring, connecting the two daughter cells, and rapidly propagates circumferentially, resulting in a hinge-like rotation.

Given that popping occurs so quickly, we questioned whether cell volume and surface area change during this process. It has been suggested that cell volume doubles at the moment of cell separation as a result of the septum expanding to cover one-half of the new spherical cell (8, 21), which would require substantial water influx over a very short time frame. To address this question, we tracked growth of individual *S. aureus* cells

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using the membrane dye FM 4-64 (Life Technologies, Grand Island, NY), and estimated cell volume and surface area from the 2D cell outlines by assuming a prolate cell shape (Fig. 2A). Overlaying the 2D cell outlines from different stages in the cell cycle (Fig. 2B, inset) revealed that a growing S. aureus cell increases both its volume and surface area throughout the cell cycle, accompanied by an overall increase in the cell aspect ratio after a small initial decrease (Fig. 2B). This small initial decrease corresponded to a phase in which the two daughter cells gradually (within minutes) became more round and more separated after popping. Following single cells and their progeny, we observed a continuous increase in cell volume for each microcolony over several generations (Fig. 2C), which is consistent with the continuous exponential volume increase that has

been described for E. coli and other bacteria (22). With respect to this continuous growth, the volume change upon popping is negligible on average (Fig. 2D). Consistent with this, using 3D deconvolution fluorescence microscopy with a S. aureus strain expressing cytoplasmic green fluorescent protein (GFP), we observed only minimal changes in cell volume and GFP intensity (fig. S3) after cell separation. Last, we estimated changes in cell surface area from 2D cell outlines by assuming a prolate cell shape, which revealed a modest net decrease in surface area upon popping (Fig. 2E and fig. S4), which is consistent with a geometric conversion from hemi-ellipsoidal to ellipsoidal shape, given constant volume. A decrease in surface area during popping indicates that the cell wall must have been under tensile stress before popping, which is in line with the hypothesis that cell-wall stress contributes to daughter-cell separation.

We next questioned whether the septum expands to become one half of the new daughter's surface (8, 2I), given that the total surface area decreases upon popping. To determine the relative contributions of the previous outer wall and septum to the surface of the new daughter, we used WGA pulse-chase labeling and 3D SIM as described above. We found that ~73% of the new daughter's surface was represented by the old wall regardless of cell size (or stage in the cell cycle) (fig. S5), similar to the ratio before cell separation (Fig. 2F), indicating that the septum constitutes only ~1/4 of the new daughter's surface and does not expand noticeably in surface area upon popping, which is contrary to the doubling of septal surface area assumed previously (8).



Fig. 1. Daughter cell separation in *S. aureus* occurs within milliseconds with characteristics of mechanical crack propagation. (A) A schematic diagram of the cell wall before daughter cell separation. (B) Snapshots of *S. aureus* strain Newman "popping" (inset) and histogram of daughter cell separation duration captured by means of phase contrast microscopy at 1000 frames/s (n = 16 popping events). (C) Distribution of cumulative counts of popping events plotted over the 2-min oscillatory period for 200 mM osmotic shocks. The red line denotes the concentration of sorbitol in the medium, and the dashed line denotes average popping counts, assuming a uniform distribution (n = 400 popping events). (D) 3D SIM images of fixed Newman cells

labeled with fluorescent WGA (WGA-488, green), which marks the outer wall and followed by 0 or 10 min of growth in the absence of WGA. Cell surfaces and septa were stained with an amine reactive dye (NHS-568, red). (**E**) Time-lapse epifluorescence images of Newman cells labeled with WGA (green) before (0 min) and after (3 min) popping. Corresponding phase-contrast (gray), membrane staining with FM 4-64 (red), and overlay of WGA and FM signals are also displayed. Two types of old wall geometry after popping were observed: hinged (left, ~80%) and nonhinged (right, ~20%). (**F**) Correlative light and SEM on Newman cells labeled with WGA followed by 10-min chase showing the two types of WGA labeling patterns as in (E). Scale bars, 1 μ m.



Fig. 2. Cell volume increases continuously throughout the cell cycle. (A) Time-lapse images of *S. aureus* cells stained with FM 4-64 (left) and outlined by fitting with ellipses (right). (B) Average aspect ratio of *S. aureus* cells throughout the cell cycle (from immediately after previous popping to ready-topop) and overlay of the cell outlines (inset) from a typical cell at different points of the cell cycle colored from blue (early) to red (late). Error bars denote standard errors (n = 27 cells). Red bars on top indicate the time fraction into the cell cycle when septation starts (left, 0.35 ± 0.03 SEM) and completes (right, 0.77 ± 0.02 SEM), respectively (n = 26 cells). (C) Representative traces of cell volume as a function of time following a microcolony starting from a single cell; solid blue traces indicate cell volumes of individual cells before popping, and the dashed black line denotes the total cell volume of all the cells

present at a given time. Cell volume and surface area were estimated from the 2D cell outlines by fitting to ellipses and assuming prolate cell shapes (that each cell was rotationally symmetric around the long axis). (**D** and **E**) Distribution of relative changes in (D) volume and (E) surface area during popping, after correcting for baseline growth rate. The black solid line represents kernel density estimate of the distribution, and the red dashed line denotes the average (2 ± 10% SD for volume, $-11 \pm 6.5\%$ SD for surface area; n = 69 division events). (**F**) (Top) 3D SIM images and corresponding extracted data (fig. S5); (bottom) fraction of old surface before (0.71 ± 0.01 SD; n = 15 cells) and after (0.73 ± 0.03 SD; n = 36 cells) popping. Cells were modeled as ellipsoids, and the contribution of the old, WGA-labeled wall to the daughter cells' total surface area was measured by fitting a plane to the old/new boundary (fig. S5A). Scale bars, 1 µm.

Additionally, the finding that the previous outer wall makes up $\sim 3/4$ of the new daughter's cell wall (as opposed to half) suggests that there must be new wall synthesis in the outer wall as well as at the septum to sustain the continuous surface expansion required for cell size homeostasis generation after generation.

Because our data suggest that accumulation of stress in the cell wall plays an important role in the ultrafast daughter cell separation, we sought to model the stress distribution in the cell wall before popping using a continuum elastostatic approach. On the basis of previous cryogenic electron microscopy data (7) and constraints determined by our experimental measurements on cell volume and surface area, we built a 3D finite element model of a "ready-to-pop" S. aureus cell wall as a prolate ellipsoidal shell with two separated septal plates that are connected with a peripheral ring (fig. S6). We assumed that the peripheral ring does not grow as much as the rest of the cell once septation begins because it is not in direct contact with the cytoplasmic membrane where new wall material is incorporated (sup-



Fig. 3. High stress in the peripheral ring prepares the cell for popping. (**A**) von Mises stress distribution in the "ready-to-pop" *S. aureus* cell wall (fig. S6, state 3) modeled as a linear elastic material (details of model construction are provided in the supplementary materials). Color represents the relative magnitude of stress. The stress at the peripheral ring (red arrow), where the cell wall splits open during popping, is higher than elsewhere in the outer wall. (**B**) Enlarged views of a cut-through slice of the cell in (A) shows high von Mises stress at the peripheral ring (red arrow) as well as the stress distribution in the circumferential and axial directions, respectively.

plementary text 1 and fig. S7). After inflating the modeled cell wall with a uniform turgor pressure in both compartments, we calculated the von Mises stress (a criterion for material failure) for the entire surface. Indeed, the von Mises stress at the peripheral ring was found to be higher than elsewhere in the outer wall (Fig. 3, A and B, and movie S6). The high von Mises stress in the



Fig. 4. Correlative SEM reveals cell cycledependent early signs of mechanical fracture. (A to D) Representative correlative fluorescence (top) and SEM images (bottom) of *S. aureus* Newman cells stained with FM 4-64, showing the cell surface features at different stages of the cell cycle. Ninety-eight percent (n = 54 cells) of the cells with visible holes [blind analysis (supplementary materials, materials and methods)] had completed septa, whereas only 49% (n = 108 cells) of the cells with completed septa had holes.

peripheral ring was mostly due to the high circumferential stress (Fig. 3B, supplementary text 2, and fig. S8) resulting from differential growth of the peripheral ring compared with the rest of the cell wall. The axial stress, which results mainly from turgor pressure (fig. S9), is actually lower in the peripheral ring than elsewhere (Fig. 3B). Modeling a series of intermediate stages throughout the growth and division cycle that includes the growth of the septum resulted in an increase in the overall cell aspect ratio during growth (movie S6), which is similar to our experimental observations (Fig. 2B).

Last, we examined the cell surface closely with high-resolution SEM in order to search for potential features of mechanical failure or fracture,

and we noticed structures in a subpopulation of cells that appeared to be perforation-like holes and cracks along the peripheral ring. Similar structures have been observed with atomic force microscopy on hydrated live cells (23). To determine whether the appearance of these structures was cell-cycle dependent, we correlated SEM with fluorescence microscopy on FM 4-64-stained cells. The holes and cracks were observed mostly on cells at later stages in the cell cycle with completed septa (Fig. 4, A to D). Specifically, 53 out of 54 cells with visible holes had completed septa (98%), whereas only 53 of 108 cells with completed septa had holes (49%), suggesting that the holes are formed after the septum is completed. Enrichment of WGA binding at the peripheral ring region was also observed when holes were present (fig. S10), suggesting that these holes are true structural changes permitting access of large proteins through the wall that are excluded at earlier stages. These perforations are likely points of mechanical failure that could initiate a propagating crack. Although the axial stress necessary for circumferential crack propagation is relatively low at the peripheral ring, the presence of perforations will lead to locally high stresses at the tips of the cracks (24) that could be sufficient to drive propagation. Consistent with this hypothesis, the distribution of the perforation lengths observed with SEM featured a cutoff length so that most perforations were <40 nm in size, which may correspond to the critical length above which cracks propagate (fig. S11A).

One candidate for creating these perforations is the enzymatic activity of cell-wall hydrolases that include Atl. Atl is the major autolysin in S. aureus and has been shown to localize at the peripheral ring region before daughter-cell separation (25, 26). In addition, atl null mutant cells form large clusters, indicating the involvement of this hydrolase in daughter cell separation (27). However, we found that the Δatl mutant still forms perforations at the peripheral ring and pops with a similar millisecond time scale as that of the wild type (fig. S12) but remains associated longer after popping, giving rise to the characteristic clustering phenotype, in which individual cells in the large clusters are nearly normal sizes and shapes. It remains possible that other cellwall hydrolases are responsible for generating the perforations we observe. However, another interesting possibility is that the perforations themselves have a mechanical origin. Given that the peripheral ring is under high circumferential stress and that the distribution of spacings between the perforations is distinctly peaked, indicating some periodicity to their locations (fig. S11B), it is plausible that perforations could form because of periodic thinning and necking of the peripheral ring under tension (supplementary text 3), a phenomenon of plastic instability observed for expanding ductile rings (28, 29).

We have described a new model for the *S. aureus* cell cycle (fig. S14) in which cell growth happens continuously throughout the cycle, and the final resolution of the peripheral ring connecting the two daughter cells happens

within a few milliseconds, with characteristics of mechanical crack propagation. It appears that this crack propagation is initiated by the perforations formed in the peripheral ring, which may result from a combination of mechanical stress and hydrolase activities, and is driven by axial wall stress originating from turgor pressure to propagate around the peripheral ring and complete daughter-cell separation. The previous outer wall of the mother cell is still partially connected immediately after cell separation, giving rise to the hingelike off-center attachment between daughter cells. This off-center attachment has been noted before (11, 30) and could explain the irregular cluster shape of Staphylococcus despite the fact that successive divisions occur in alternating perpendicular planes (11, 31, 32). The mechanical crack propagation-driven model for S. aureus daughter-cell separation described here demonstrates how a biological system with appropriate geometrical design can use mechanical forces to achieve large-scale morphological changes.

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SUPPLEMENTARY MATERIALS

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www.sciencemag.org/content/348/6234/574/suppl/DC1 Materials and Methods Supplementary Text Figs. S1 to S14 Table S1 References (*32–36*) Movies S1 to S7 27 October 2014; accepted 25 March 2015

PROTEIN DYNAMICS

Direct observation of hierarchical protein dynamics

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One of the fundamental challenges of physical biology is to understand the relationship between protein dynamics and function. At physiological temperatures, functional motions arise from the complex interplay of thermal motions of proteins and their environments. Here, we determine the hierarchy in the protein conformational energy landscape that underlies these motions, based on a series of temperature-dependent magic-angle spinning multinuclear nuclear-magnetic-resonance relaxation measurements in a hydrated nanocrystalline protein. The results support strong coupling between protein and solvent dynamics above 160 kelvin, with fast solvent motions, slow protein side-chain motions, and fast protein backbone motions being activated consecutively. Low activation energy, small-amplitude local motions dominate at low temperatures, with larger-amplitude, anisotropic, and functionally relevant motions involving entire peptide units becoming dominant at temperatures above 220 kelvin.

roteins must traverse complex conformational energy landscapes to perform their physiological function. This is achieved through thermally activated fluctuations (1), and when specific molecular motions cease at low temperatures, function also ceases or becomes reduced (2, 3). Understanding the hierarchy of these motions thus holds the key to understanding how proteins function on a molecular level at physiological temperatures. Contrary to expectations, and despite large differences in structure and function between proteins of different families, dynamic properties appear to exhibit common general features, in particular apparent transitions between different dynamic regimes as a function of temperature. These transitions are thought to occur as a result of coupling between proteins and the surrounding solvent (4).

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*Corresponding author. E-mail: j.r.lewandowski@warwick.ac.uk (J.R.L.); martin.blackledge@ibs.fr (M.B.); lyndon.emsley@epfl.ch (L.E.) †Present address: Department of Chemistry, University of Warwick, Coventry CV4 7AL, UK. ‡Present address: Department of Chemistry, University of York, York YO10 5DD, UK. Observing and understanding protein dynamic transitions has been a focus of many fields of research over the past 40 years, including neutron scattering (5), Mössbauer spectroscopy (4, 6), Terahertz spectroscopy (7), dielectric spectroscopy (4), differential scanning calorimetry (8), x-ray crystallography (2, 9), and molecular dynamics simulation (10, 11). This relative wealth of information has not, however, led to a consensus picture of dynamical transitions and their origins, with different techniques detecting distinct processes, leading to apparently contradictory descriptions (12, 13). This may be due to the widely varying conditions required for the diverse techniques or to the sensitivity of the different physical measurements to dynamics occurring on different time scales.

Here we measure, in a single sample, a set of 13 different nuclear magnetic resonance (NMR) observables that are sensitive to dynamics occurring on different time scales and in different parts of the system over temperatures from 105 to 280 K in the fully hydrated crystalline protein GB1, a small globular protein specifically binding to antibodies. The analysis of multiple probes that report on the different structural components of this complex system allows us to develop a complete and coherent picture of the dynamic processes across the whole temperature range, as





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